

Conformationally Dependent Photodynamics of Glycine

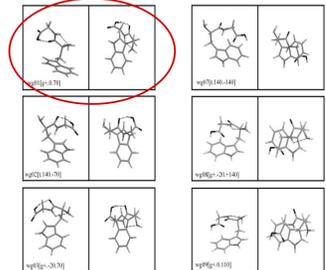
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Motivation



1 H. Valdes et al. *J. Phys. Chem. B* 110 (2006) 6385.

• Trp-Gly-Gly

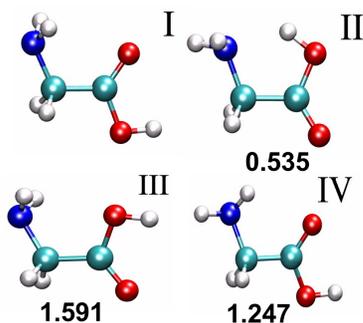
the most populated structure in the gas phase is not observed experimentally (REMPI and UV-UV and IR-UV hole burning spectroscopy)

→ possible explanation is that very short lifetime of some conformers prevents them from being measured¹

The difference among calculated conformers is due to various H bonding pattern in Gly-Gly part → glycine represents suitable starting point

• Addressing a general question about the possible role of ground state conformations on excited state dynamics.

Glycine in the ground state



Glycine is a nonrigid (floppy) system → a large number of conformers

The minima are separated by small barriers that can be easily overcome (highest barrier is ~ 940 (III) and 740 (IV) cm⁻¹ for interconversion to global minimum I)

4 energetically lowest structures in the gas phase plus relative energies in kcal/mol²

2 J. Chocholoušová et al. *J. Chem. Phys.* 113 (2000) 4629.

Vertical excitation energies (oscillator strength in parenthesis) for conformer 4.

	CASPT2 6-31G**	CASPT2 ANO-L ³	EOM-CCSD aD(T) ⁴	experiment ⁵
	5.76 eV	5.65 eV (0.001)	5.88 eV (0.0005)	~5.8-6.0 eV
	6.68 eV	6.09 eV (0.005)	6.32 eV (0.0124)	

the vertical excitation energies are similar for all conformers, the ordering of states is the same.

3 L.-S. Andres *J. Am. Chem. Soc.* 118 (1996) 12200. 4 A. Osted et al *J. Phys. Chem. A* 109 (2005) 1430. 5 Inagaki, T. *Biopolymers* 12 (1973) 1353.

Possible photophysical pathways

Localization of conical intersection CASSCF/6-31G** SA3(6,8), in cases of H-atom loss the s and p diffuse functions were added (with exponent 0.02)

The energies and the character of excited states was checked by using the RIC2 method with cc-pVTZ and aug-cc-pVDZ on N and O basis set with added diffuse s and p functions on H atom (with coefficient 0.02) on the geometries of the conical intersections

Dynamic correlation included by using CASPT2 method (level shift 0.3 Hartree)

S3/S2

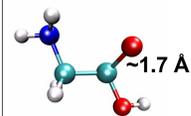
For all conformers the S3/S2 conical intersection is connected with NH₂ group planarization and moderate elongation of N-H bond (from 1.099 Å for conformer 4 to 1.210 Å for conformer 2, 3).

The character of S3 state is changed to n_N → π* when slightly moving the hydrogen atom → the basis set used is not excellent but should be sufficient

S2/S1

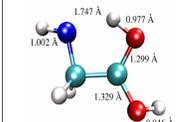
• C=O bond stretching – found for almost all conformers

• the deexcitation process leads to **vibrationally excited molecule** in the ground state.



• H-atom transfer – found only for conformer 4

• the energy of the CI is lower than in the case of C=O bond stretching (~ 1.1 eV)
• **question** – the stability of dihydroxyisomer in the ground state.

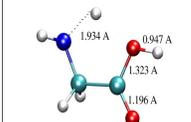


the H-atom transfer was observed in the case of β-turn studied by A. L. Sobolewski

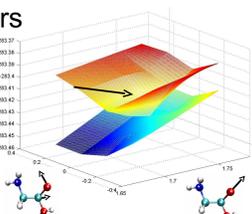
• H-atom loss – found for conformers 2 and 3

• energy of the CI is approximately the same as in the case of C=O bond stretching
• the channel leads to a **fragmented molecule** in the ground state

the dissociation along the N-H coordinate was observed for in neutral biological aromatic molecules experimentally as well as described theoretically⁶



3 A. L. Sobolewski, W. Domcke et al *Phys. Chem. Chem. Phys.* 4 (2002) 1093.



ground and excited state surfaces (CASPT2 energies) in respect to g and h vector that can be projected to C=O bond stretching and -CO₂ bending

Ab initio non adiabatic molecular dynamics Full Multiple Spawning⁷

Time-dependent Schroedinger equation is solved

Wave function ansatz $\Psi = \sum_I \chi_I(R; t) \phi_I(r; R)$

time-dependent χ_I
nuclear wavefunction associated with Ith state

electronic wavefunction ϕ_I

Nuclear wavefunction – superposition of frozen Gaussian basis functions

$\chi_I(R; t) = \sum_{j=1}^{N_I(t)} C_j^I(t) \chi_j^I(R, R_j^I(t), P_j^I(t), \gamma_j^I(t), \alpha_j^I)$

time-dependent coefficients $C_j^I(t)$

Individual, multidimensional nuclear basis functions - one-dimensional Gaussian basis functions:

$e^{-\alpha_j^I(t) \prod_{\rho=1}^{3N} \left(\frac{2\alpha_{\rho j}^I}{\Pi} \right)^{1/4} \exp(-\alpha_{\rho j}^I(t) (R_{\rho j} - R_{\rho j}^I(t))^2 + i P_{\rho j}^I(t) (R_{\rho j} - R_{\rho j}^I(t)))}$

nuclear phase $\alpha_j^I(t)$ propagated semi classically
time-independent width $\frac{2\alpha_{\rho j}^I}{\Pi}$
position and momentum propagated classically

Substituting the wavefunction ansatz into nuclear Schrodinger equation - evolution of the coefficients

$$\frac{dC^I(t)}{dt} = -i(S_{II}^{-1}) \left((H_{II} - i\dot{S}_{II})C^I + \sum H_{IJ}C^J \right)$$

The new basis functions are “spawned” in region of high non-adiabatic coupling

7 M. Ben-Nun, J. Quenneville, T. J. Martinez *J. Phys. Chem. A* 104 (2000) 5171.

Simulation details

Electronic structure

CASSCF wavefunction – SA3(4,4) with 6-31G** basis set

Dynamics

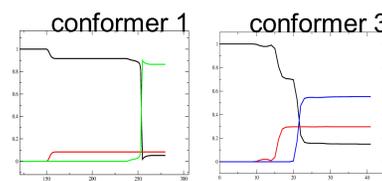
Initial trajectories sampled according to Wigner distribution function computed for the initial vibrational ground state

Timestep 10 a.u. (0.24 fs)

Dynamics started on S3 using for each conformer its ground state minimum geometry

Results – work in progress

S3/S2



The AIMS data show almost complete population transfer from S3 to S2 in tens of femtoseconds.

Conformers 1 and 4 behave equally they are interconverting faster that the population transfer to S2 takes place → they are **indistinguishable**

S2/S1 - observed channels

Channels observed - C=O bond stretching and H-atom transfer (only for conformers 1 and 4). The H-atom loss channel was not found due to small basis set used in the simulations (there was only extensive N-H bond stretching (~1.2 Å)).

No rotation around C-C bond was observed before the S2/S1 interconversion takes place → conformers 1,4 and 2,3 are **distinguishable** in respect to H-atom transfer.

Outlook

the basis set used is not suitable for all photophysical pathways (does not allow H-atom loss). More trajectories also need to be calculated to get relevant of information about photophysics in the S2 state.

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